

# AN UNSTEADY PULSATILE FLOW OF SOME MHD NON-NEWTONIAN NANOFLUIDS WITH HALL CURRENT AND ION SLIP THROUGH A POROUS MEDIUM

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## ABSTRACT

*The current study examines the pulsatile flow of MHD Non-Newtonian nanofluids with hall current, ion slip and thermal radiation. Using regular perturbation technique, velocity and unsteady temperature distributions for both Eyring-Powell nanofluid and Casson nano fluid are analysed by suitable transformations. When we compare Eyring-Powell nanofluid model with Casson nano fluid model it is ob that u is high for Casson nanofluid model, and Eyring-Powell nano fluid shows the highest temperature variations.*

**KEYWORDS:** Pulsatile Flow, Eyring-Powell Nanofluid, Casson Nanofluid, Hall Current & Ion Slip and Thermal Radiation

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## INTRODUCTION

The study of Non-Newtonian fluids has a vital role in fields like food processing, polymer processing, etc. Hayat et al. [1] Nadeem and Saleem[2], Tasawar Hayat et al. [3], Macharla Jayachandra Babu et al. [4], Vijayalakshmi and Shankar Bandari. [5], Asha and Sunitha[6], Mallick and Misra [7] and Misra and Sinha [8] are some of the researchers who have focused on MHD Eyring-Powell Nanofluid in recent years. A few recent investigators that may be mentioned are Motsa and Shateyi[9], Shehzad et al.[10], Rizwanulhaq et al. [11], Mustafa and Junaid Ahmad Khan [12], Imran Ullah et al. [13], Sreekala and Kesavareddy [14], Kamran et al. [15], Rajib Biswas and Sardar Firoz Ahmmed [16], Pudhari srilatha[17], Jannath Begam et al.[18], Deivanayaki et al.[19], Jannath Begam et al. [20]and Jannath Begam et al. [21].

Here, we attempt to examine the hall current and ion slip effects on pulsatile flow of some Non-Newtonian fluids with thermal radiation.

## MATHEMATICAL FORMULATION

The constitutive equation for Eyring-Powell nanofluid and Casson nanofluid is

$$\tau_{ij} = \mu_{nf} \frac{\partial v_i}{\partial x_j} + \frac{1}{\beta} \sin h^{-1} \left( \frac{1}{c} \frac{\partial v_i}{\partial x_j} \right) \quad (1)$$

Since  $\sin h^{-1} x \approx x$ ,  $|x| \ll 1$ , then

$$\begin{aligned}\tau_{ij} &= \mu_{nf} \frac{\partial V_i}{\partial x_j} + \frac{1}{\beta} \left( \frac{1}{c} \frac{\partial V_i}{\partial x_j} \right) \\ &= \mu_{nf} \left( 1 + \frac{1}{\beta c \mu_{nf}} \right) \frac{\partial V_i}{\partial x_j}\end{aligned}\quad (2)$$

$$\tau_{ij} = \begin{cases} 2 \left( \mu B + \frac{P_{y^*}}{\sqrt{2\pi}} \right) e_{ij}, \pi > \pi_c \\ 2 \left( \mu B + \frac{P_{y^*}}{\sqrt{2\pi_c}} \right) e_{ij}, \pi < \pi_c \end{cases}$$

$$\tau_{ij} = \mu_{nf} \left( 1 + \frac{1}{\beta} \right) \frac{\partial V_i}{\partial x_j} \quad (3)$$

Unsteady incompressible Eyring-Powell nanofluid and Casson nanofluid are examined, with its x-axis selected as lower plate and transversely applied magnetic field along y axis. Using Boussinesq approximation, we have

$$\frac{\partial u^*}{\partial t^*} = -\frac{1}{\rho_{nf}} \frac{\partial P^*}{\partial x^*} + q \vartheta_{nf} \left( 1 + \frac{1}{\beta c \mu_{nf}} \right) \frac{\partial^2 u^*}{\partial y^{*2}} + (1-q) q \vartheta_{nf} \left( 1 + \frac{1}{\beta} \right) \frac{\partial^2 u^*}{\partial y^{*2}} - \frac{\mu_{nf}}{\rho_{nf}} u^* + \frac{\sigma_{nf} B_0^2 (1 + \beta_e \beta_i)}{\rho_{nf} [(1 + \beta_e \beta_i)^2 + \beta_e^2]} u^* \quad (4)$$

$$\frac{\partial T^*}{\partial t^*} = \frac{k_{nf}}{(\rho C_P)_{nf}} \frac{\partial^2 T^*}{\partial y^{*2}} + \frac{1}{(\rho C_P)_{nf}} \frac{\partial q_r}{\partial y^*} + q \frac{\mu_{nf} \left( 1 + \frac{1}{\beta c \mu_{nf}} \right)}{(\rho C_P)_{nf}} \left( \frac{\partial u^*}{\partial y^*} \right)^2 + (1-q) \frac{\mu_{nf} \left( 1 + \frac{1}{\beta} \right)}{(\rho C_P)_{nf}} \left( \frac{\partial u^*}{\partial y^*} \right)^2 + \frac{Q_0}{(\rho C_P)_{nf}} (T^* - T_0) \quad (5)$$

In equations (4) and (5),  $q \in [0, 1]$ , it states Eyring-Powell nanofluid for  $q=1$  and Casson nanofluid for  $q=0$ .

and

$$u^* = 0, T^* = T_0 \text{ at } y^* = 0 \quad (6)$$

$$u^* = 0, T^* = T_1 \text{ at } y^* = h \quad (7)$$

The physical properties of nanofluid such as  $\mu_{nf}$ ,  $\rho_{nf}$ ,  $(\rho C_P)_{nf}$  and  $k_{nf}$  are given as

$$\mu_{nf} = \frac{\mu_f}{(1-\phi)^{2.5}} \quad (8)$$

$$\rho_{nf} = (1-\phi)\rho_f + \phi\rho_s \quad (9)$$

$$(\rho C_P)_{nf} = (1-\phi)(\rho C_P)_f + \phi(\rho C_P)_s \quad (10)$$

$$\frac{k_{nf}}{k_f} = \frac{k_s + 2k_f - 2\phi(k_f - k_s)}{k_s + 2k_f + \phi(k_f - k_s)} \quad (11)$$

$$q_r = -\frac{4\sigma^*}{3k^*} \frac{\partial T^{*4}}{\partial y^*} \quad (12)$$

$$T^{*4} \cong 4T_1^3 T^* - 3T_1^4 \quad (13)$$

Using Equations (12) and (13), Equation (5) becomes

$$\frac{\partial T^*}{\partial t^*} = \frac{k_{nf}}{(\rho C_P)_{nf}} \frac{\partial^2 T^*}{\partial y^{*2}} + \frac{1}{(\rho C_P)_{nf}} \left[ \frac{16\sigma^*}{3k^*} T_1^3 \frac{\partial^2 T^*}{\partial y^{*2}} \right] + q \frac{\mu_{nf} \left( 1 + \frac{1}{\beta c \mu_{nf}} \right)}{(\rho C_P)_{nf}} \left( \frac{\partial u^*}{\partial y^*} \right)^2 + (1-q) \frac{\mu_{nf} \left( 1 + \frac{1}{\beta} \right)}{(\rho C_P)_{nf}} \left( \frac{\partial u^*}{\partial y^*} \right)^2 + \frac{Q_0}{(\rho C_P)_{nf}} (T^* - T_0) \quad (14)$$

$$-\frac{1}{\rho_f} \frac{\partial P^*}{\partial x^*} = A(1 + \epsilon e^{i\omega t}) \quad (15)$$

Introducing the following dimensionless variables

$$u = \frac{u^* \omega}{A}, t = t^* \omega, x = \frac{x^*}{h}, y = \frac{y^*}{h}, \theta = \frac{T^* - T_0}{T_1 - T_0}, P = \frac{P^*}{A_p \mu_f h}, k_0 = \frac{1}{\beta c \mu_f} \quad (16)$$

And parameters as in Equations (8) to (11) into Equations (15), (4) and (14), we get

$$-\frac{\partial P}{\partial x} = 1 + \epsilon e^{it} \quad (17)$$

$$\frac{\partial u}{\partial t} = -\frac{1}{A_1} \frac{\partial P}{\partial x} + q \frac{A_2}{A_1} \left(1 + \frac{k_0}{A_2}\right) \frac{1}{R} \frac{\partial^2 u}{\partial y^2} + (1 - q) \frac{A_2}{A_1} \left(1 + \frac{1}{\beta}\right) \frac{1}{R} \frac{\partial^2 u}{\partial y^2} - \frac{A_2}{A_1} \left(\frac{1}{RD_a} + M_1\right) u \quad (18)$$

$$\frac{\partial \theta}{\partial t} = \left(\frac{A_4}{A_3} + \frac{4Rd}{3A_3}\right) \frac{1}{RP_r} \frac{\partial^2 \theta}{\partial y^2} + q \frac{A_2}{A_3} \left(1 + \frac{k_0}{A_2}\right) \frac{Ec}{R} \left(\frac{\partial u}{\partial y}\right)^2 + (1 - q) \frac{A_2}{A_3} \left(1 + \frac{1}{\beta}\right) \frac{Ec}{R} \left(\frac{\partial u}{\partial y}\right)^2 + \frac{Q}{A_3 R} \theta \quad (19)$$

and

$$u = 0, \theta = 0 \text{ at } y = 0 \quad (20)$$

$$u = 0, \theta = 1 \text{ at } y = 1 \quad (21)$$

$$\text{where } A_1 = (1 - \phi) + \phi \frac{\rho_s}{\rho_f}, A_2 = \frac{1}{(1 - \phi)^{2.5}},$$

$$A_3 = (1 - \phi) + \phi \frac{(\rho C_P)_s}{(\rho C_P)_f}, A_4 = \frac{k_s + 2k_f - 2\phi(k_f - k_s)}{k_s + 2k_f + \phi(k_f - k_s)},$$

$$Pr = \frac{(\rho C_P)_f \nu_f}{k_f}, Ec = \frac{A^2}{\omega^2 (\rho C_P)_f (T_1 - T_0)}, R = \frac{\omega h^2}{\mu_f}, k_0 = \frac{1}{\beta c \mu_f}, M = \frac{\sigma_{nf} B_0^2}{\rho_{nf}}$$

$$Da = \frac{k}{h^2}, Rd = \frac{4\sigma^* T_1^3}{k_f k^*}, Q = \frac{Q_0 h^2}{(\rho C_P)_f \mu_f}, M_1 = \frac{M(1 + \beta_e \beta_i)}{(1 + \beta_e \beta_i)^2 + \beta_e^2}$$

## SOLUTION FOR THE PROBLEM

The  $u$  and  $\theta$  are written as

$$u = u_0(y) + \epsilon u_1(y) e^{it} \quad (22)$$

$$\theta = \theta_0(y) + \epsilon \theta_1(y) e^{it} + \epsilon^2 \theta_2(y) e^{2it} \quad (23)$$

Using Equations (17), (22) and (23) in Equations (18) and (19) and comparing the terms of like powers of  $\epsilon$ , we get

$$q \frac{A_2}{A_1} \left(1 + \frac{k_0}{A_2}\right) \frac{1}{R} u_0'' + (1 - q) \frac{A_2}{A_1} \left(1 + \frac{1}{\beta}\right) \frac{1}{R} u_0'' - \frac{A_2}{A_1} \left(\frac{1}{RD_a} + M_1\right) u_0 = -\frac{1}{A_1} \quad (24)$$

$$q \frac{A_2}{A_1} \left(1 + \frac{k_0}{A_2}\right) \frac{1}{R} u_1'' + (1 - q) \frac{A_2}{A_1} \left(1 + \frac{1}{\beta}\right) \frac{1}{R} u_1'' - \frac{A_2}{A_1} \left(\frac{1}{RD_a} + M_1\right) u_1 - i u_1 = -\frac{1}{A_1} \quad (25)$$

$$\left(\frac{A_4}{A_3} + \frac{4Rd}{3A_3}\right) \frac{1}{RP_r} \theta_0'' + q \frac{A_2}{A_3} \left(1 + \frac{k_0}{A_2}\right) (u_0')^2 + (1 - q) \frac{A_2}{A_3} \left(1 + \frac{1}{\beta}\right) (u_0')^2 + \frac{Q}{A_3 R} \theta_0 = 0 \quad (26)$$

$$\left(\frac{A_4}{A_3} + \frac{4Rd}{3A_3}\right) \frac{1}{RP_r} \theta_1'' + 2q \frac{A_2}{A_3} \left(1 + \frac{k_0}{A_2}\right) u_0' u_1' + 2(1 - q) \frac{A_2}{A_3} \left(1 + \frac{1}{\beta}\right) u_0' u_1' \frac{Q}{A_3 R} \theta_1 - i \theta_1 = 0 \quad (27)$$

$$\left(\frac{A_4}{A_3} + \frac{4Rd}{3A_3}\right) \frac{1}{RP_r} \theta_2'' + 2q \frac{A_2}{A_3} \left(1 + \frac{k_0}{A_2}\right) (u_1')^2 + 2(1 - q) \frac{A_2}{A_3} \left(1 + \frac{1}{\beta}\right) (u_1')^2 + \frac{Q}{A_3 R} \theta_2 - 2i \theta_2 = 0 \quad (28)$$

and

$$u_0 = 0, \theta_0 = 0 \text{ at } y = 0 \quad (29)$$

$$= 0, \theta_0 = 1 \text{ at } y = 1 \quad (30)$$

$$u_1 = 0, \theta_1 = 0, \theta_2 = 0 \text{ at } y = 0 \quad (31)$$

$$u_1 = 0, \theta_1 = 0, \theta_2 = 0 \text{ at } y = 1 \quad (32)$$

By solving equations (24) to (28) with the corresponding boundary conditions (29) to (32), we obtain

$$u_0 = B_7 e^{\sqrt{B_3} y} + B_6 e^{-\sqrt{B_3} y} - \frac{B_4}{B_3} \quad (33)$$

$$u_1 = B_9 e^{\sqrt{B_5} y} + B_8 e^{-\sqrt{B_5} y} - \frac{B_4}{B_5} \quad (34)$$

$$\theta_0 = c_{11} \cos \sqrt{d_2} y + c_{12} \sin \sqrt{d_2} y + c_8 e^{2\sqrt{B_3} y} + c_9 e^{-2\sqrt{B_3} y} + c_{10} \quad (35)$$

$$\theta_1 = D_6 e^{\sqrt{C_4} y} + D_5 e^{-\sqrt{C_4} y} + D_1 e^{m_1 y} + D_2 e^{m_2 y} + D_3 e^{m_3 y} + D_4 e^{m_4 y} \quad (36)$$

$$\theta_2 = D_{11} e^{\sqrt{C_6} y} + D_{10} e^{-\sqrt{C_6} y} + D_7 e^{2\sqrt{B_5} y} + D_8 e^{-2\sqrt{B_5} y} + D_9 \quad (37)$$

## RESULTS AND DISCUSSIONS

To find the physical insight, the velocity and temperature profiles are studied for different parameters and discussed graphically. Figure (1a) pictures, velocity diminishes for  $k_0$  and the opposite behavior has keenly been observed from figure (1b) for the Casson nanofluid parameter for different  $\beta_e$ . Velocity diminishes for ion slip parameter  $\beta_i$  as illustrated in figure (2a). The contrasting effect has been pointed out in figure (2b). Velocity diminishes as  $k_0$  and  $\beta$  increase as examined in figure (3a) and (3b). Casson nanofluid shows little high increase in unsteady temperature for various  $Ec$  than  $k_0$  as analysed in figure (4).  $\beta$  shows slightly increasing  $\theta_t$  than  $k_0$  for various heat source parameter  $Q$  is observed from Figure (5).

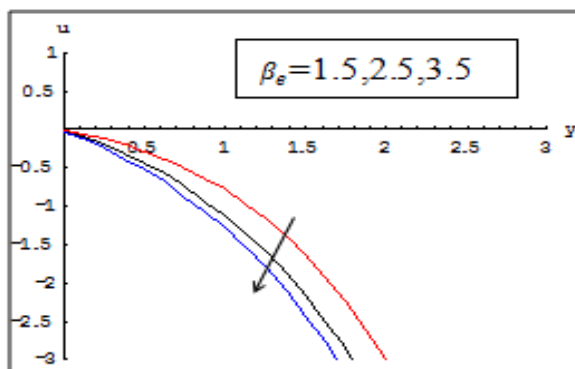


Figure (1a): Impact of  $\beta_e$  on Velocity Profile  $u$  of Eyring-Powell Nanofluid

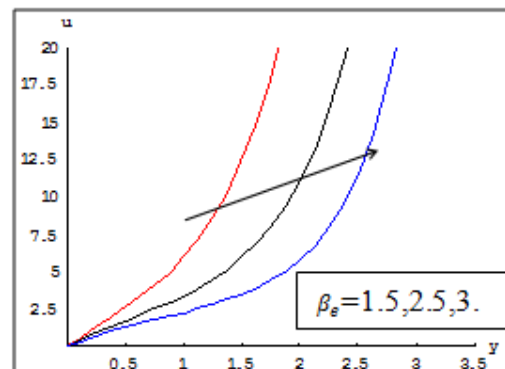


Figure (1b): Impact of  $\beta_e$  on Velocity Profile  $u$  Cass on Nanofluid

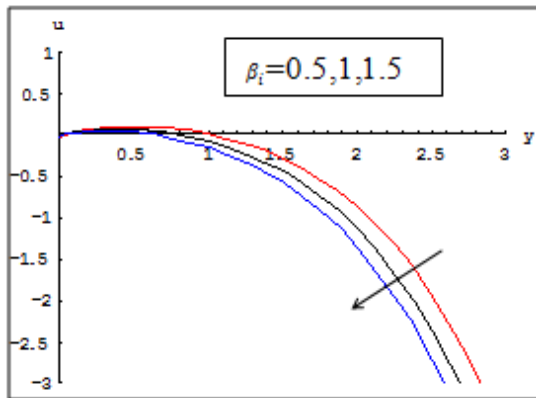


Figure (2a): Impact of  $\beta_i$  on Velocity Profile  $u$  of Eyring-Powell Nanofluid

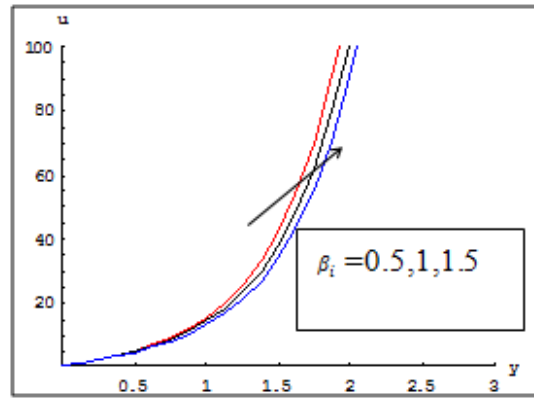


Figure (2b): Impact of  $\beta_i$  on Velocity Profile  $u$  of Cass on Nanofluid

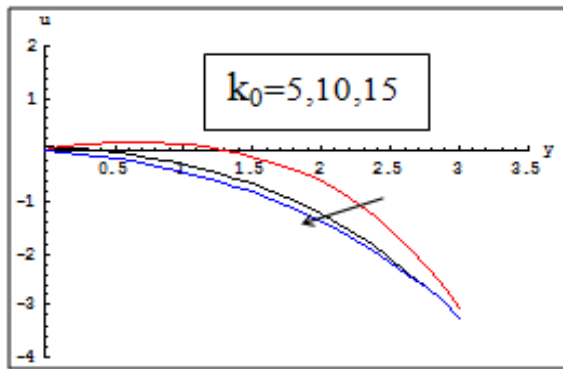


Figure (3a): Impact of  $k_0$  on Velocity Profile  $u$  of Eyring-Powell Nanofluid

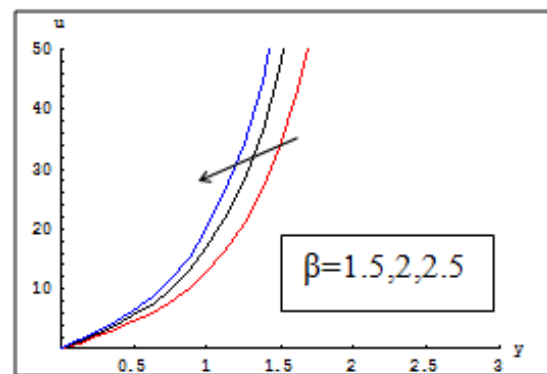


Figure (3b): Impact of  $\beta$  on Velocity Profile  $u$  of Cass on Nanofluid

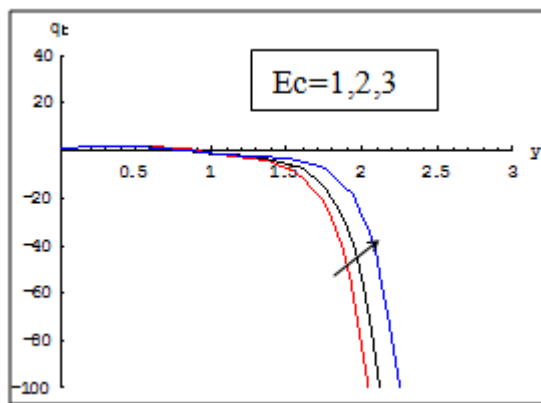


Figure (4a): Impact of  $Ec$  on Unsteady Temperature Profile  $\theta_t$  of Eyring-Powell Nanofluid

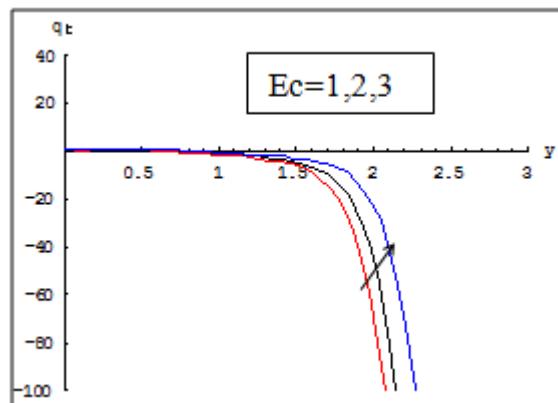
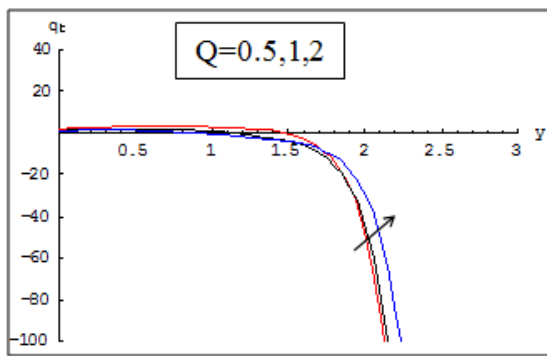
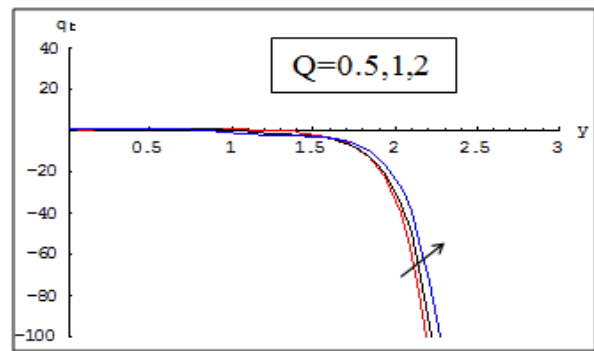


Figure (4b): Impact of  $Ec$  on Unsteady Temperature Profile  $\theta_t$  of Cass on Nanofluid



**Figure (5a): Impact of Q on Unsteady Temperature Profile  $\theta_t$  of Eyring-Powell Nanofluid**



**Figure (5b): Impact of Q on Unsteady Temperature Profile  $\theta_t$  of Cass on Nanofluid**

## CONCLUSIONS

Finally, we conclude that,

- The velocity profile decreases with Hall current and ion slip effect for  $k_0$  and increase for  $\beta$ .
- The velocity profile gets decreased with increasing values of  $k_0$  and  $\beta$ , and the temperature profile gets increased with enhanced  $Ec$  and  $Q$  for both  $k_0$  and  $\beta$ .

Although a very slight difference is observed between  $k_0$  and  $\beta$ , we conclude that  $k_0$  shows minute variations when compared with  $\beta$ .

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## APPENDIX

### NOMENCLATURE

$\beta$	Casson nanofluid parameter
$Q_0$	Coefficient of Heat source/Sink
Da	Darcy Number
$\rho_{nf}$	Density of Nanofluid
$\rho_f$	Density of the Base fluid
$\rho_s$	Density of the Nanoparticle
h	Distance Between the plates
$\mu_{nf}$	Dynamic viscosity of the Nanofluid
Ec	Eckert Number
$(\rho C_p)_{nf}$	Effective Heat Capacitance of Nanofluid
$k_0$	Eyring-Powell nanofluid parameter
R	Frequency parameter
$\beta_e$	Hall current parameter
$(\rho C_p)_f$	Heat Capacitance of Base fluid
$(\rho C_p)_s$	Heat Capacitance of Nanoparticle
Q	Heat Source parameter
$\beta_i$	Ion slip parameter
$k^*$	Mean Absorption Coefficient
$\phi$	Nanoparticles volume fraction
k	Permeability of porous medium
Pr	Prandtl number
P*	Pressure
Rd	Radiation Parameter
$q_r$	Radiative Heat flux
$\sigma^*$	Stefan-Boltzmann Constant
$k_f$	Thermal conductivity of Base fluid
$k_{nf}$	Thermal conductivity of Nanofluid
$k_s$	Thermal conductivity of Nanoparticle
$u^*$	Velocity component in x*-direction
$\mu_f$	Viscosity of the Base fluid